

7.1-1 Eigenfunction expansion -- 1

Consider
$$\begin{cases} Lu \equiv \frac{d}{dx} \left[p(x) \frac{du}{dx} \right] + q(x)u = f(x), & a < x < b \\ B_1 u = u_0 \\ B_2 u = v_0 \end{cases}$$

**regular Sturm-Liouville
eigenvalue problem**

Associated eigenvalue problem

$$\begin{cases} L\phi(x) + \lambda\rho(x)\phi(x) = 0, & a < x < b \\ B_1\phi = 0 \\ B_2\phi = 0 \end{cases}$$

C^1 : continuously differentiable
 C^0 : continuous

Let (i) $p(x) > 0$ & $p(x) \in \underline{C^1}$, $q(x), \rho(x) \in \underline{C^0}$ and

(ii) B_1 & B_2 be ■ unmixed B. C. or
■ periodic B. C., with $p(a) = p(b)$

Charles Sturm (1803-1855) and Joseph Liouville (1809-1882) worked in 1836, 1837 for the above boundary value problem.



$$\frac{d}{dx} \left[p(x) \frac{d\phi(x)}{dx} \right] + q(x)\phi(x) + \lambda \rho(x)\phi(x) = 0, \quad \begin{array}{l} p(x) > 0 \text{ \& } p(x) \in C^1 \\ q(x), \rho(x) \in C^0 \end{array}$$

(self-adjoint)

(1) there exists a **countable** set of eigenvalues, which form an increasing sequence: $\lambda_1 < \lambda_2 < \dots < \lambda_n < \dots \rightarrow \infty$

real, simple, countable eigenvalues

$$\begin{cases} \underline{L\phi(x) + \lambda\rho(x)\phi(x)} = 0 \\ B_1\phi = 0 \\ B_2\phi = 0 \end{cases} \quad \begin{array}{c} \text{take complex} \\ \text{conjugate} \Rightarrow \end{array} \begin{cases} \underline{L\bar{\phi}(x) + \bar{\lambda}\rho(x)\bar{\phi}(x)} = 0 \\ B_1\bar{\phi} = 0 \\ B_2\bar{\phi} = 0 \end{cases} \quad \begin{array}{c} (\bar{\lambda}, \bar{\phi}) \text{ is also} \\ \text{an eigen-pair} \end{array}$$

$$\begin{aligned} & \int_a^b \left\{ \bar{\phi} [\underline{L\phi + \lambda\rho\phi}] - \phi [\underline{L\bar{\phi} + \bar{\lambda}\rho\bar{\phi}}] \right\} dx = 0 \\ & = \int_a^b \left\{ \bar{\phi} \lambda \rho \phi - \phi \bar{\lambda} \rho \bar{\phi} \right\} dx + \cancel{J[\phi, \bar{\phi}] \Big|_a^b} \quad \text{(Self-adjoint!)} \\ & = (\lambda - \bar{\lambda}) \int_a^b \phi \rho \bar{\phi} dx \quad \Rightarrow \quad \lambda = \bar{\lambda} \end{aligned}$$

$$\frac{d}{dx} \left[p(x) \frac{d\phi(x)}{dx} \right] + q(x)\phi(x) + \lambda \rho(x)\phi(x) = 0$$

(2) If $\phi_i(x)$ & $\phi_j(x)$ correspond to λ_i & λ_j , and $\lambda_i \neq \lambda_j$ then

$$\int_a^b \rho(x) \phi_i(x) \phi_j(x) dx = 0$$

$$\begin{cases} L\phi_i(x) + \lambda_i \rho(x)\phi_i(x) = 0 \\ B_1\phi_i = 0 \\ B_2\phi_i = 0 \end{cases}$$

$$\begin{cases} L\phi_j(x) + \lambda_j \rho(x)\phi_j(x) = 0 \\ B_1\phi_j = 0 \\ B_2\phi_j = 0 \end{cases}$$

$$\int_a^b \left\{ \phi_j [L\phi_i + \lambda_i \rho \phi_i] - \phi_i [L\phi_j + \lambda_j \rho \phi_j] \right\} dx = 0$$

$$= \int_a^b \left\{ \phi_j \lambda_i \rho \phi_i - \phi_i \lambda_j \rho \phi_j \right\} dx + J[\phi_i, \phi_j] \Big|_a^b \quad (\text{Self-adjoint!})$$

$$= (\lambda_i - \lambda_j) \int_a^b \phi_i \rho \phi_j dx \quad \Rightarrow \quad \int_a^b \rho(x) \phi_i(x) \phi_j(x) dx = 0$$

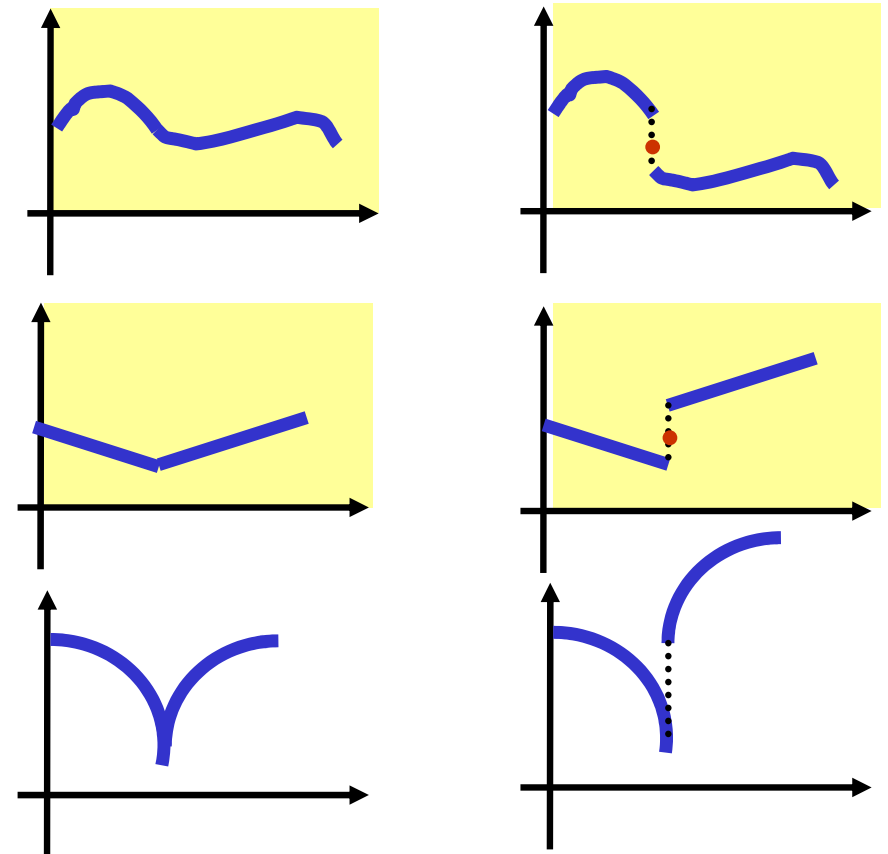
$$\frac{d}{dx} \left[p(x) \frac{d\phi(x)}{dx} \right] + q(x)\phi(x) + \lambda \rho(x)\phi(x) = 0$$

(3) Eigenfunctions $\{\phi_n(x)\}$ forms a complete set in $[a, b]$, then piecewise smooth function $u(x)$ can be expressed as

$$u(x) = \sum_{n=1}^{\infty} c_n \phi_n(x)$$

where $c_n = \int_a^b \rho(x) u(x) \phi_n(x) dx$

with $\int_a^b \rho \phi_n^2 dx = 1$



7.1-5 Remarks -- 1

$$\begin{cases} Lu \equiv \frac{d}{dx} \left[p(x) \frac{d\phi(x)}{dx} \right] + q(x)\phi(x) + \lambda \rho(x)\phi(x) = 0 & , a < x < b \\ B_1 \phi = 0 \\ B_2 \phi = 0 \end{cases}$$

Regular Sturm-Liouville eigenvalue problem

- (i) $p(x) > 0$ & $p(x) \in C^1$
 $q(x), \rho(x) \in C^0$ and
- (ii) B_1 & B_2 be
 - unmixed B. C. or
 - periodic B. C., with $p(a) = p(b)$

Singular Sturm-Liouville eigenvalue problem

- (i) $p(x) = 0$ at some points in $[a, b]$; or
- (ii) $q(x)$ is discontinuous in $[a, b]$
- (iii) $[a, b]$ is a semi-infinite or an infinite interval

7.1-6 Remarks -- 2

Consider $Lu \equiv a_2(x)u'' + a_1(x)u' + a_0(x)u = f(x)$, $a < x < b$

If $a_2'(x) \neq a_1(x)$, and $a_2(x) \neq 0$ in $[a, b]$

i.e. L is not formally self-adjoint

To make the equation self-adjoint, find a $r(x)$ such that :

$$r(x)[a_2(x)u'' + a_1(x)u' + a_0(x)u] = \frac{d}{dx} \left[p(x) \frac{du}{dx} \right] + q(x)u \quad ??$$

$$\Rightarrow \text{want } \begin{cases} r(x)a_2(x) = p(x) \\ r(x)a_1(x) = p'(x) \end{cases}$$

$$\Rightarrow \frac{p'(x)}{p(x)} = \frac{a_1(x)}{a_2(x)}$$

$$\Rightarrow p'(x) = \frac{a_1(x)}{a_2(x)} p(x)$$

$$\frac{dp}{p} = \frac{a_1(x)}{a_2(x)} dx$$

$$\therefore p(x) = \exp \left[\int \frac{a_1(x)}{a_2(x)} dx \right]$$

$$r(x) = \frac{1}{a_2(x)} \exp \left[\int \frac{a_1(x)}{a_2(x)} dx \right]$$

rLu is self-adjoint !

7.1-7 Remarks -- 3

regular Sturm-Liouville eigenvalue problem

$$\begin{cases} \frac{d}{dx} \left[p \frac{d\phi}{dx} \right] + q\phi + \lambda\rho\phi = 0, & a < x < b \\ B_1\phi = 0 \\ B_2\phi = 0 \end{cases}$$

self-adjoint problem

- real eigenvalues
- eigenfunctions are orthogonal
- eigenfunctions is a complete set

If the eigen-problem is not formally self-adjoint, then the eigen-functions of the problem will not form an orthogonal set. But it could form a bi-orthogonal set with the eigen-function of the adjoint problem.

7.1-8 Remarks -- 4

Consider

(adjoint eigenvalue problem)

$$\begin{cases} \underline{L\phi + \lambda\rho\phi} = 0, & a < x < b \\ B_1\phi = 0 \\ B_2\phi = 0 \end{cases} \quad \begin{cases} \underline{L^*\psi + \lambda^*\rho\psi} = 0, & a < x < b \\ B_1^*\psi = 0 \\ B_2^*\psi = 0 \end{cases}$$

$$\int_a^b \left\{ \psi \left[\underline{L\phi + \lambda\rho\phi} \right] - \phi \left[\underline{L^*\psi + \lambda^*\rho\psi} \right] \right\} dx = 0$$

$$= \int_a^b \left\{ \psi \lambda \rho \phi - \phi \lambda^* \rho \psi \right\} dx + \cancel{J[\phi, \psi] \Big|_a^b}$$

$$= (\lambda - \lambda^*) \int_a^b \phi \rho \psi dx \quad \Rightarrow \quad \int_a^b \rho(x) \phi(x) \psi(x) dx = 0$$

($\phi(x)$ and $\psi(x)$ are orthogonal)

7.2-1 BVP (with homogeneous B.C.)

$$(I) \quad \begin{cases} Lu = f(x) , a < x < b \\ u(a) = 0 \\ u(b) = 0 \end{cases}$$

$$(II) \quad \begin{cases} Lu_H = 0 , a < x < b \\ u_H(a) = 0 \\ u_H(b) = 0 \end{cases}$$

$$(III) \quad \begin{cases} L^* u_H^* = 0 , a < x < b \\ u_H^*(a) = 0 \\ u_H^*(b) = 0 \end{cases}$$

$$(IV) \quad \begin{cases} LG = \delta(x - \xi) \\ G(a; \xi) = 0 \\ G(b; \xi) = 0 \end{cases}$$

7.2-1 BVP (with homogeneous B.C.)

$$(I) \quad \begin{cases} \underline{Lu = f(x)} , a < x < b \\ u(a) = 0 \\ u(b) = 0 \end{cases} \quad \begin{cases} \underline{L\phi_n(x) + \lambda_n \rho(x)\phi_n(x) = 0} , a < x < b \\ \phi_n(a) = 0 \\ \phi_n(b) = 0 \end{cases} \quad \int_a^b \rho(x)\phi_m(x)\phi_n(x)dx = \delta_{mn}$$

$$u(x) = \sum_{n=1}^{\infty} c_n \phi_n(x) \quad \text{where} \quad c_n = \int_a^b \rho(x) u(x) \phi_n(x) dx$$

$$\int_a^b \left\{ \phi_n [Lu - f] - u [L\phi_n + \lambda_n \rho \phi_n] \right\} dx = 0$$

$$c_n = \frac{\int_a^b \rho(x) u(x) \phi_n(x) dx}{\int_a^b \rho(\eta) \phi_n^2(\eta) d\eta}$$

$$= \int_a^b \left\{ -\phi_n f - u \lambda_n \rho \phi_n \right\} dx + J[u, \phi_n] \Big|_a^b$$

$$c_n = \frac{\int_a^b \phi_n(x) f(x) dx}{-\lambda_n}$$

$$u(x) = \sum_{n=1}^{\infty} \left[\frac{\int_a^b \phi_n(\eta) f(\eta) d\eta}{-\lambda_n} \right] \phi_n(x)$$

7.2-2 Example 1--1

$$\begin{cases} u'' = f(x), & 0 < x < \ell \\ u(0) = 0 \\ u(\ell) = 0 \end{cases}$$

$$\begin{cases} \phi_n''(x) + \lambda_n \phi_n(x) = 0, & 0 < x < \ell \\ \phi_n(0) = 0 \\ \phi_n(\ell) = 0 \end{cases}$$

eigenvalues &
eigenvectors

$$\lambda_n = (n\pi)^2 / \ell^2$$

$$\phi_n(x) = \sin \frac{n\pi x}{\ell}$$

Method 1:

$$u(x) = \sum_{n=1}^{\infty} c_n \sin \frac{n\pi x}{\ell}$$

from D.E.

$$-\sum_{n=1}^{\infty} c_n \left(\frac{n\pi}{\ell}\right)^2 \sin \frac{n\pi x}{\ell} = f(x)$$

$$-c_n \cdot \left(\frac{n\pi}{\ell}\right)^2 \cdot \frac{\ell}{2} = \int_0^{\ell} f(x) \sin \frac{n\pi x}{\ell} dx$$

$$u(x) = \frac{-2\ell}{\pi^2} \sum_{n=1}^{\infty} \left[\frac{\int_0^{\ell} f(\eta) \sin \frac{n\pi \eta}{\ell} d\eta}{n^2} \right] \sin \frac{n\pi x}{\ell}$$

$$\begin{cases} \underline{u'' = f(x)} , & 0 < x < \ell \\ u(0) = 0 \\ u(\ell) = 0 \end{cases} \quad \begin{cases} \underline{\phi_n''(x) + \lambda_n \phi_n(x) = 0} , & 0 < x < \ell \\ \phi_n(0) = 0 \\ \phi_n(\ell) = 0 \end{cases}$$

Method 2:

$$u(x) = \sum_{n=1}^{\infty} c_n \sin \frac{n\pi x}{\ell} \quad c_n = \frac{2}{\ell} \int_0^{\ell} u(\eta) \sin \frac{n\pi \eta}{\ell} d\eta$$

$$\int_0^{\ell} \{ \phi_n [\underline{u'' - f}] - u [\underline{\phi_n'' + \lambda_n \rho \phi_n}] \} dx = 0$$

$$= \int_0^{\ell} \{ -\phi_n f - u \lambda_n \phi_n \} dx + \cancel{J[u, \phi_n] \Big|_0^{\ell}}$$

$$\therefore -c_n \cdot \left(\frac{n\pi}{\ell}\right)^2 \cdot \frac{\ell}{2} = \int_0^{\ell} f(x) \sin \frac{n\pi x}{\ell} dx$$

$$u(x) = \dots$$

$$\begin{cases} u'' = f(x) , & 0 < x < \ell \\ u(0) = 0 \\ u(\ell) = 0 \end{cases}$$

$$\begin{cases} u(x) = \sum_{n=1}^{\infty} c_n \sin \frac{n\pi x}{\ell} \\ c_n = \frac{2}{\ell} \int_0^{\ell} u(x) \sin \frac{n\pi x}{\ell} dx \end{cases}$$

Method 3:

$$\begin{cases} c_n = \frac{2}{\ell} \int_0^{\ell} u(x) \sin \frac{n\pi x}{\ell} dx \\ u(x) = \sum_{n=1}^{\infty} c_n \sin \frac{n\pi x}{\ell} \end{cases}$$

finite transform !!
c.f. Laplace or Fourier
transform

Apply transform $\frac{2}{\ell} \int_0^{\ell} (*) \sin \frac{n\pi x}{\ell} dx$ to the D.E.

$$\int_0^{\ell} u''(x) \sin \frac{n\pi x}{\ell} dx = \int_0^{\ell} f(x) \sin \frac{n\pi x}{\ell} dx$$

Integration by
parts + B.C.

$$-c_n \cdot \left(\frac{n\pi}{\ell}\right)^2 \cdot \frac{\ell}{2} = \int_0^{\ell} f(x) \sin \frac{n\pi x}{\ell} dx$$

$$u(x) = \dots$$

7.3-1 BVP (with inhomogeneous B.C.)

$$\begin{cases} u'' = f(x) , & 0 < x < \ell \\ u(0) = u_0 \\ u(\ell) = v_0 \end{cases}$$

$$\begin{cases} \phi_n''(x) + \lambda_n \phi_n(x) = 0 , & 0 < x < \ell \\ \phi_n(0) = 0 \\ \phi_n(\ell) = 0 \end{cases}$$

eigenvalues &
eigenvectors

$$\lambda_n = (n\pi)^2 / \ell^2$$

$$\phi_n(x) = \sin \frac{n\pi x}{\ell}$$

Method 1:

$$u(x) = \sum_{n=1}^{\infty} d_n \sin \frac{n\pi x}{\ell}$$

from D.E.

$$-\sum_{n=1}^{\infty} d_n \left(\frac{n\pi}{\ell}\right)^2 \sin \frac{n\pi x}{\ell} = f(x)$$

$$-d_n \cdot \left(\frac{n\pi}{\ell}\right)^2 \cdot \frac{\ell}{2} = \int_0^{\ell} f(x) \sin \frac{n\pi x}{\ell} dx$$

$$u(x) = \frac{-2\ell}{\pi^2} \sum_{n=1}^{\infty} \left[\frac{\int_0^{\ell} f(\eta) \sin \frac{n\pi \eta}{\ell} d\eta}{n^2} \right] \sin \frac{n\pi x}{\ell}$$

????

$$\begin{cases} \underline{u'' = f(x)} , & 0 < x < \ell \\ u(0) = u_0 \\ u(\ell) = v_0 \end{cases} \quad \begin{cases} \underline{\phi_n''(x) + \lambda_n \phi_n(x) = 0} , & 0 < x < \ell \\ \phi_n(0) = 0 \\ \phi_n(\ell) = 0 \end{cases}$$

Method 2: $u(x) = \sum_{n=1}^{\infty} d_n \sin \frac{n\pi x}{\ell} \quad d_n = \frac{2}{\ell} \int_0^{\ell} u(\eta) \sin \frac{n\pi \eta}{\ell} d\eta$

$$\int_0^{\ell} \{ \phi_n [\underline{u'' - f}] - u [\underline{\phi_n'' + \lambda_n \phi_n}] \} dx = 0$$

$$= \int_0^{\ell} \{ -\phi_n f - u \lambda_n \phi_n \} dx - (u \phi_n' - \phi_n u') \Big|_0^{\ell}$$

$$\therefore -d_n \cdot \left(\frac{n\pi}{\ell} \right)^2 \cdot \frac{\ell}{2} = \int_0^{\ell} f(x) \sin \frac{n\pi x}{\ell} dx - \frac{n\pi}{\ell} \left[\underline{u_0 - (-1)^n v_0} \right]$$

$$d_n = -\frac{2\ell}{\pi^2} \frac{1}{n^2} \int_0^{\ell} f(x) \sin \frac{n\pi x}{\ell} dx + \frac{2}{n\pi} \left[u_0 - (-1)^n v_0 \right]$$

$$\begin{cases} u'' = f(x) , & 0 < x < \ell \\ u(0) = u_0 \\ u(\ell) = v_0 \end{cases}$$

7.3-3 Example 2--3

Method 3: Apply transform $\frac{2}{\ell} \int_0^\ell (*) \sin \frac{n\pi x}{\ell} dx$ to the D.E.

$$\int_0^\ell u''(x) \sin \frac{n\pi x}{\ell} dx = \int_0^\ell f(x) \sin \frac{n\pi x}{\ell} dx$$

Integration by parts

$$-c_n \cdot \left(\frac{n\pi}{\ell}\right)^2 \cdot \frac{\ell}{2} + \frac{n\pi}{\ell} [u_0 - (-1)^n v_0] = \int_0^\ell f(x) \sin \frac{n\pi x}{\ell} dx$$

$$u(x) = \dots$$

$$\begin{cases} u'' = f(x) , & 0 < x < \ell \\ u(0) = u_0 \\ u(\ell) = v_0 \end{cases}$$

7.3-4 Example 2--4

Method 4: Let $u(x) = U(x) + u_0 \left(1 - \frac{x}{\ell}\right) + v_0 \frac{x}{\ell}$

$$\begin{cases} U'' = f(x) , & 0 < x < \ell \\ U(0) = 0 \\ U(\ell) = 0 \end{cases}$$

Converge faster (see lecture notes)



$$\therefore u(x) = u_0 \left(1 - \frac{x}{\ell}\right) + v_0 \frac{x}{\ell} + \frac{-2\ell}{\pi^2} \sum_{n=1}^{\infty} \left[\frac{\int_0^{\ell} f(\eta) \sin \frac{n\pi\eta}{\ell} d\eta}{n^2} \right] \sin \frac{n\pi x}{\ell}$$

From method 2:

$$u(x) = \frac{-2\ell}{\pi^2} \sum_{n=1}^{\infty} \left[\frac{\int_0^{\ell} f(\eta) \sin(n\pi\eta/l) d\eta}{n^2} \right] \sin \frac{n\pi x}{l} + \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{1}{n} \left[u_0 - (-1)^n v_0 \right] \sin \frac{n\pi x}{l}$$

7.4-1 Example 3--1 (not-self-adjoint)

$$\begin{cases} u'' + 4u' + 3u = 3, & 0 < x < 1 \\ u(0) = 0 \\ u(1) = 0 \end{cases}$$

Method 1:

$$u = u_h + u_p = A e^{-x} + B e^{-3x} + 1$$

By using of B.C., one has.....

Method 2: Green's function approach

$$u = 3 \int_0^1 G(x; \xi) d\xi$$

$$\begin{cases} u'' + 4u' + 3u = 3, & 0 < x < 1 \\ u(0) = 0 \\ u(1) = 0 \end{cases}$$

Example 3-2 (not-self-adjoint)

Method 3:

Associated eigenvalue problem

$$\begin{cases} \phi'' + 4\phi' + \lambda\phi = 0, & 0 < x < 1 \\ \phi(0) = 0 \\ \phi(1) = 0 \end{cases}$$

eigenvalues &
eigenvectors

$$\lambda_n = (n\pi)^2 + 4$$

$$\phi_n(x) = e^{-2x} \sin n\pi x$$

$$u(x) = e^{-2x} \sum_{n=1}^{\infty} A_n \sin n\pi x$$

from D.E.

$$e^{-2x} \sum_{n=1}^{\infty} A_n (-n^2\pi^2 - 1) \sin n\pi x = 3$$

**Orthogonality of
eigenfunctions ??**

$$A_n = ???$$

$$\begin{cases} u'' + 4u' + 3u = 3, & 0 < x < 1 \\ u(0) = 0 \\ u(1) = 0 \end{cases}$$

Example 3-3 (not-self-adjoint)

Associated eigenvalue problem

$$\begin{cases} \phi'' + 4\phi' + \lambda\phi = 0, & 0 < x < 1 \\ \phi(0) = 0 \\ \phi(1) = 0 \end{cases}$$

$$\int_0^1 \phi_n(x) \phi_m(x) dx \neq 0$$

Adjoint eigenvalue problem

$$\begin{cases} \psi'' - 4\psi' + \lambda\psi = 0, & 0 < x < 1 \\ \psi(0) = 0 \\ \psi(1) = 0 \end{cases}$$

$$\phi_n(x) = e^{-2x} \sin n\pi x$$

$$\lambda_n = (n\pi)^2 + 4$$

$$\psi_n(x) = e^{2x} \sin n\pi x$$

$$e^{-2x} \sum_{n=1}^{\infty} A_n (-n^2\pi^2 - 1) \sin n\pi x = 3$$

From orthogonality of eigenfunctions

$$A_n \cdot (-n^2\pi^2 - 1) \cdot \frac{1}{2} = 3 \int_0^1 e^{2x} \sin n\pi x dx \quad \therefore A_n = \dots$$

$$\begin{cases} u'' + 4u' + 3u = 3, & 0 < x < 1 \\ u(0) = 0 \\ u(1) = 0 \end{cases}$$

Associated eigenvalue problem

$$\begin{cases} \phi'' + 4\phi' + \lambda\phi = 0, & 0 < x < 1 \\ \phi(0) = 0 \\ \phi(1) = 0 \end{cases}$$

Method 4:

Associated eigenvalue problem

$$\begin{cases} e^{4x}\phi'' + 4e^{4x}\phi' + \lambda e^{4x}\phi = 0, & 0 < x < 1 \\ \phi(0) = 0 \\ \phi(1) = 0 \end{cases}$$

(e^{4x} = integrating factor)

$$\lambda_n = (n\pi)^2 + 4$$

$$\phi_n(x) = e^{-2x} \sin n\pi x$$

$$u(x) = e^{-2x} \sum_{n=1}^{\infty} A_n \sin n\pi x$$

from D.E.

$$e^{2x} \sum_{n=1}^{\infty} A_n (-n^2\pi^2 - 1) \sin n\pi x = 3$$

**Orthogonality of
eigenfunctions OK!!**

7.4-5 Example 3--5 (not-self-adjoint)

Method 5:

Assume $u(x) = e^{-2x} \sum_{n=1}^{\infty} A_n \sin n\pi x = e^{-2x} w(x)$

$$\begin{cases} w'' - w = 3e^{2x}, & 0 < x < 1 \\ w(0) = 0 \\ w(1) = 0 \end{cases} \quad \text{(Self-adjoint now!)}$$

$$w(x) = \sum_{n=1}^{\infty} B_n \sin n\pi x$$

from D.E. ...